

Newton's Method and Rate of Convergence.

Inverting $DP(x_n)$ at each step is tantamount to solving a linearized version of the equation $P(x_n) = 0$. This is computationally expensive. There is however a gain — faster (here quadratic) convergence.

Suppose Newton's method converges to a solution $x^* \in X$ of the equation $P(x) = 0$, and $(DP(x^*))^{-1}$ exists. Furthermore, assume that within an open region containing x^* and the sequence $\{x_n\}$, the quantities $\|DP(x)\|$, $k=2,3$ and $\|(DP(x))^{-1}\|$ are bounded.

Then, letting $T(x) = x - (DP(x))^{-1}P(x)$,

$$\begin{aligned} x_{n+1} - x^* &= x_n - (DP(x_n))^{-1}P(x_n) - x^* \\ &= x_n - (DP(x_n))^{-1}P(x_n) \\ &\quad - (x^* - (DP(x^*))^{-1}P(x^*)) \end{aligned}$$

(by hypothesis $P(x^*) = 0$)

$$= T(x_n) - T(x^*)$$

$$\begin{aligned} DT(x) &= (DP(x^*))^{-1} D^2P(x^*) (DP(x^*))^{-1} P(x^*) \\ &= 0 \end{aligned}$$

By proposition 2 (Lecture Notes 9(b)), on mean values,

$$\|x_{n+1} - x^*\| \leq \frac{1}{2} \sup_{\bar{x}} \|D^2 T(\bar{x})\| \|x_n - x^*\|^2$$

where $\bar{x} = x_n + \alpha(x_n - x^*)$, $0 \leq \alpha \leq 1$.

Have $\|x_{n+1} - x^*\| \leq c \|x_n - x^*\|^2$

where $c = \frac{1}{2} \sup_{x \in R} \|D^2 T(x)\|$, a bound depending upon $\|D^3 P(x)\|$.

This is quadratic convergence.

Application 1

Euler Lagrange equations in the calculus of variations can be cast in the form

$$(EL) \quad \dot{x} = F(x(t))$$

$$C x(t_1) = c_1$$

$$D x(t_2) = d_2$$

for problems with fixed end points. Such a boundary value problem can be solved by Newton's method. First (EL) can be written as a

$$x(t) = x(t_1) + \int_{t_1}^t F(x(s)) ds$$

$$t_1 \leq t \leq t_2$$

$\Leftrightarrow P(x) = 0$ where,

$$(P(x))(t) = x(t) - x(t_1) - \int_{t_1}^t F(x(\sigma)) d\sigma$$

$t_1 \leq t \leq t_2$

It is left as an exercise to show that the Fréchet derivative $DP(x)$ is given by

$$(DP(x) \cdot h)(t) = h(t) - \int_{t_1}^t F_x(x(\sigma)) h(\sigma) d\sigma$$

where $F_x(x(\sigma)) = \left[\frac{\partial F^i}{\partial x_j} \right]_{x(\sigma)}$.

Newton's method is then equivalent to

$$DP(x_n) \cdot (x_{n+1} - x_n) = -P(x_n)$$

\Leftrightarrow

$$\begin{aligned} x_{n+1}(t) - x_n(t) &= \int_{t_1}^t F_x(x_n(\sigma)) (x_{n+1}(\sigma) - x_n(\sigma)) d\sigma \\ &= -x_n(t) + x(t_1) + \int_{t_1}^t F(x_n(\sigma)) d\sigma \end{aligned}$$

Differentiating both sides we get,

$$\begin{aligned} \dot{x}_{n+1}(t) &= F_x(x_n(t)) \cdot (x_{n+1}(t) - x_n(t)) + F(x_n(t)) \\ &= F_x(x_n(t)) \cdot x_{n+1}(t) + \left(F(x_n(t)) - F_x(x_n(t)) x_n(t) \right) \end{aligned}$$

which is a linear time varying system with input,

of the form $\dot{x}_{n+1}(t) = A_n(t) x_{n+1}(t) + b_n(t)$

where

$$A_n(t) = F_x(x_n(t))$$

$$b_n(t) = F(x_n(t)) - F_x(x_n(t))x_n(t)$$

add to this the boundary conditions

$$C x_{n+1}(t_1) = c_1$$

$$D x_{n+1}(t_2) = d_2$$

We can solve this linear boundary value problem as follows.

Let $\bar{\Phi}_n(t_2, t)$ be defined by the transition matrix equation

$$\dot{\bar{\Phi}}_n(t_2, t) = -\bar{\Phi}_n(t_2, t) A_n(t)$$

$$\bar{\Phi}_n(t_2, t_2) = \underline{1}$$

Then ~~the~~

$$x_{n+1}(t_2) = \bar{\Phi}_n(t_2, t_1) x(t_1) + \int_{t_1}^{t_2} \bar{\Phi}_n(t_2, \sigma) b_n(\sigma) d\sigma$$

The boundary conditions become

$$C x_{n+1}(t_1) = c_1$$

$$D \bar{\Phi}_n(t_2, t_1) x_{n+1}(t_1) = d_2 - D \int_{t_1}^{t_2} \bar{\Phi}_n(t_2, \sigma) b_n(\sigma) d\sigma$$

Solve for $x_{n+1}(t)$. It then follows that

$$x_{n+1}(t) = \underline{\Phi}_n(t, t_1) x(t_1) + \int_{t_1}^t \underline{\Phi}_n(t, \sigma) b_n(\sigma) d\sigma$$

where

$\underline{\Phi}_n(t, t_1)$ satisfies

$$\dot{\underline{\Phi}}_n(t, t_1) = A_n(t) \underline{\Phi}_n(t, t_1)$$

$$\underline{\Phi}_n(t_1, t_1) = \mathbb{1}.$$

Of course

$$\underline{\Phi}_n(t, t_1) = \underline{\Phi}_n(t, t_2) \underline{\Phi}_n(t_2, t_1)$$

$$= \underline{\Phi}_n(t_2, t)^{-1} \underline{\Phi}_n(t_2, t_1)$$

Thus we solve a linear system at each step of the Newton iteration scheme.

This is a practical method to solve Euler Lagrange equations.